

SCATTERING FOR THE GROSS-PITAEVSKII EQUATION

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ABSTRACT. We investigate the asymptotic behavior at time infinity of solutions close to a non-zero constant equilibrium for the Gross-Pitaevskii (or Ginzburg-Landau-Schrödinger) equation. We prove that, in dimensions larger than 3, small perturbations can be approximated at time infinity by the linearized evolution, and the wave operators are homeomorphic around 0 in certain Sobolev spaces.

1. Introduction

The cubic Schrödinger equation for $\varphi(t, x) : \mathbb{R} \times \mathbb{R}^d \rightarrow \mathbb{C}$,

$$(1.1) \quad i\partial_t \varphi = -\Delta \varphi + |\varphi|^2 \varphi,$$

has numerous physical applications. Noting that the single mode uniform oscillation e^{-it} is trivially a solution, and writing $\varphi = e^{-it}\psi$, we arrive at

$$(1.2) \quad i\partial_t \psi = -\Delta \psi + (|\psi|^2 - 1)\psi, \quad \psi(0, \cdot) = \psi_0,$$

the *Ginzburg-Landau-Schrödinger* or *Gross-Pitaevskii* equation, which models the dynamics of Bose-Einstein condensates, superfluids, and superconductors (see [7] for a review). In these physical settings, $\varphi = e^{-it}$ (or $\psi = 1$) corresponds to a stationary, constant-density condensate. Typically, one is interested in local perturbations of such an equilibrium, in which case the natural boundary condition as $|x| \rightarrow \infty$ should *not* be the condition $\psi \rightarrow 0$ usually imposed when studying (1.1), but rather

$$(1.3) \quad |\psi| \rightarrow 1 \quad \text{as} \quad |x| \rightarrow \infty.$$

Equation (1.2) with boundary condition (1.3) possesses, in dimension $d = 2$, well-known static, spatially localized, topologically non-trivial *vortex* solutions of the form $\psi(x) = e^{in \cdot \arg(x)} f(|x|)$, $n \in \mathbb{Z}$. There are also traveling-wave solutions, $\psi(t, x) = v(x_1 - ct, x_2, \dots, x_d)$, with v resembling a pair of vortices ($d = 2$), or a ring of vorticity ($d \geq 3$): [3, 2, 4, 11]. For various Ginzburg-Landau Schrödinger and wave-type equations (such as (1.2)), the dynamics of vortex points ($d = 2$), filaments ($d = 3$), etc., has been extensively studied in certain asymptotic limits: see, e.g., [6, 22, 16, 21, 17, 27, 13] and references therein. (In addition, there is a large literature on vortex dynamics in parabolic versions of (1.2).)

Very little, however, seems to be known about the long-time behaviour of solutions of (1.2)–(1.3), and, in particular, about the stability of the various interesting vortex-type solutions mentioned above (however see, eg., [18, 23] for

numerical and formal asymptotic studies of vortex traveling wave stability, and [12] for vortex stability/instability in the gauge-invariant wave-equation version of (1.2)).

As a first step in this direction, we consider here a much simpler problem: asymptotic stability of the *vacuum* solution

$$(1.4) \quad \psi(t, x) \equiv e^{i\theta_0} = \text{constant} \in S^1.$$

By the gauge change $\psi \rightarrow e^{-i\theta_0}\psi$, we may assume $\theta_0 = 0$. Further writing $\psi = 1 + u$, we have

$$(1.5) \quad \begin{aligned} i\partial_t u + \Delta u - 2\operatorname{Re} u &= F(u), \\ F(u) &:= (u + 2\bar{u} + |u|^2)u. \end{aligned}$$

The main goal of this paper is to investigate the asymptotic stability of the vacuum in terms of the scattering theory around the constant solution (1.4) of (1.2), or the 0 solution of (1.5).

The first step is to introduce the following change of variable in order to represent the free (linear) evolution as a unitary group:

$$(1.6) \quad u \mapsto v := V^{-1}u := U^{-1}\operatorname{Re} u + i\operatorname{Im} u, \quad U := \sqrt{-\Delta(2 - \Delta)}^{-1}.$$

Then v satisfies the equation

$$(1.7) \quad i\partial_t v - \sqrt{-\Delta(2 - \Delta)}v = -iV^{-1}iF(Vv).$$

The singularity of U^{-1} at the Fourier origin seems a natural effect of the interaction with the nonzero constant background in (1.2). Due to this singularity, together with a quadratic nonlinearity which is not gauge invariant, the scattering problem for (1.5) turns out to be much more difficult than the original NLS (1.1) with spatially decaying data (i.e. $\varphi \in L^2$). Thus we have obtained so far the scattering only for higher dimensions $d \geq 4$. Global wellposedness of (1.5) in the energy space $u \in H^1(\mathbb{R}^d)$ was proved in [3] for $d = 2, 3$. It is based on a priori bounds of the energy and so does not work for $d \geq 4$, where the cubic nonlinearity is either energy-critical or super critical. Our theorem gives global wellposedness for $d \geq 4$ and small data as a by-product.

Theorem 1.1 (Stability and scattering). *Suppose that $d \geq 4$ and $|\sigma| \leq \frac{d-3}{2} - \frac{1}{d}$. If $U^\sigma V^{-1}u(0)$ is sufficiently small in $H^{d/2-1}(\mathbb{R}^d)$, then $U^\sigma V^{-1}u(t)$ remains small in $H^{d/2-1}(\mathbb{R}^d)$ for all $t \in \mathbb{R}$. Moreover, there exist $v_\pm \in U^{-\sigma}H^{d/2-1}(\mathbb{R}^d)$ such that*

$$(1.8) \quad \left\| U^\sigma(V^{-1}u(t) - e^{-iHt}v_\pm) \right\|_{H^{d/2-1}(\mathbb{R}^d)} \rightarrow 0$$

as $t \rightarrow \pm\infty$, where $H = \sqrt{-\Delta(2 - \Delta)}$, and the wave operators $v_\pm \mapsto V^{-1}u(0)$ are local homeomorphisms around 0 in $U^{-\sigma}H^{d/2-1}(\mathbb{R}^d)$.

The corresponding result for the cubic NLS (1.1) holds [5] for any $d \geq 2$ and $|\sigma| \leq d/2 - 1$. Moreover, the wave operators are defined on the whole energy space H^1 for $d = 2, 3, 4$, and are bijective for $d = 3, 4$ [10, 25]. For $d = 1$, modified

wave operators have been constructed in [24] for small data in weighted Sobolev spaces.

The same result for NLS with the nonlinearity $F(u)$ is easily obtained for $d \geq 4$ and $|\sigma| \leq d/2 - 2$ by the argument in [5]. The lower dimensional case is more difficult due to the quadratic terms and slow decay of the free evolution. For $d = 3$, the wave operators have been constructed together with their inverses for the nonlinearities $|u|u$, u^2 and \bar{u}^2 for small data in weighted Sobolev spaces [9, 14]. For $d = 2$, the modified wave operators for $|u|u$ and the ordinary ones for u^2 and \bar{u}^2 have been constructed for small data in weighted Sobolev spaces [8, 15]. It has been proved that $|u|^2$ in $d = 2$ does not admit the ordinary wave operators [26].

The estimates in weighted spaces are quite sensitive to the linear part and so do not easily extend to our problem. Here the difficulty in our equation is that the Galilean invariance is also destroyed by the boundary condition. For the Maxwell-Higgs equation, a similar hyperbolic problem, the Lorentz invariance is preserved and the scattering result has been obtained in [28].

To prove the above theorem, we first derive L^p decay and Strichartz estimates for the linear evolution e^{-iHt} . These are slightly better for $d \geq 3$ around the Fourier origin than the free Schrödinger equation, because of curvature properties of the characteristic surfaces. This is the reason why we have a wider range of σ than the NLS with the same nonlinearity. Actually, this gain is strong enough for $d \geq 5$ to cancel the U^{-1} singularity in (1.7), but not enough for $d \leq 4$. Instead, we use a seemingly miraculous change of the unknown function – a kind of “low frequency normal form” transformation – which completely removes the singularities at the Fourier origin from the nonlinearity. However, it still leaves regular quadratic nonlinearities, whose decay is out of our control for $d \leq 3$.

In the next section 2, we investigate the linearized operator, deriving the decay and Strichartz estimates for the evolution. In Section 3, we perform the change of unknown function, and then in Section 4, prove the main theorem.

2. Linearized operator

In this section, we derive the L^p decay and Strichartz estimate for the linear equation

$$(2.1) \quad i\dot{u} + \Delta u - 2\operatorname{Re} u = f.$$

As in the introduction, we define the following \mathbb{C} -linear Fourier multipliers

$$(2.2) \quad U := \sqrt{-\Delta(2 - \Delta)^{-1}}, \quad H := \sqrt{-\Delta(2 - \Delta)},$$

and \mathbb{R} -linear operator

$$(2.3) \quad Vu := U \operatorname{Re} u + i \operatorname{Im} u.$$

By using the identities $H = -\Delta U^{-1} = (2 - \Delta)U$, we obtain

$$(2.4) \quad i\dot{u} + \Delta u - 2\operatorname{Re} u = iV(\partial_t + iH)V^{-1}u.$$

Thus the linear evolution by (2.1) is reduced to the unitary group e^{-iHt} , for which we have the following decay and Strichartz estimates. Here $\dot{B}_p^s := \dot{B}_{p,2}^s$ denotes the homogeneous Besov spaces (cf. [1]), and we denote spacetime norms by $\|f\|_{L^p X} := \left(\int_{-\infty}^{\infty} \|f(t, \cdot)\|_X^p dt\right)^{1/p}$, etc.

Theorem 2.1. *Let $d \geq 1$. (i) For any $2 \leq q \leq \infty$, we have*

$$(2.5) \quad \|e^{-itH}\varphi\|_{\dot{B}_q^0} \lesssim t^{-d\sigma} \|U^{(d-2)\sigma}\varphi\|_{\dot{B}_{q'}^0},$$

where $q' = q/(q - 1)$ is the dual exponent and $\sigma = 1/2 - 1/q$.

(ii) For $j = 1, 2$, let $2 \leq p_j, q_j \leq \infty$, $2/p_j + d/q_j = d/2$ and $s_j = \frac{d-2}{2}(1/2 - 1/q_j)$, but $(p_j, q_j) \neq (2, \infty)$. Then we have

$$(2.6) \quad \begin{aligned} & \|e^{-itH}\varphi\|_{L^{p_1}\dot{B}_{q_1}^0} \lesssim \|U^{s_1}\varphi\|_{L^2}, \\ & \left\| \int_{-\infty}^t e^{-i(t-s)H} f(s) ds \right\|_{L^{p_1}\dot{B}_{q_1}^0} \lesssim \|U^{s_1+s_2} f\|_{L^{p_2'}\dot{B}_{q_2'}^0} \end{aligned}$$

When $d > 2$, these estimates are slightly better than those for the free Schrödinger equation because of the factor U , which is small around the Fourier origin. This gain is due to the fact that the curvature of the characteristic surface $|\xi|\sqrt{2 + |\xi|^2}$ is larger than that of $|\xi|^2$ in the $(d - 1)$ angular directions and smaller in the radial directions. The above theorem follows from the standard argument for the stationary phase. For completeness, we give a proof in the following subsection.

2.1. Stationary phase estimate.

Theorem 2.2. *Let $\varphi(r) \in C^\infty(0, \infty)$ satisfy the following.*

1. $\varphi'(r), \varphi''(r) > 0$ for all $r > 0$.
2. $\varphi'(r) \sim \varphi'(s)$ and $\varphi''(r) \sim \varphi''(s)$ for $0 < s < r < 2s$.
3. $|\varphi^{(1+k)}(r)| \lesssim \varphi'(r)/r^k$ for all $r > 0$ and $k \in \mathbb{N}$.

Let $\chi(r)$ be a dyadic cut-off function which is supported around $r \sim R$ and satisfies

$$(2.7) \quad |\chi^{(k)}(r)| \lesssim R^{-k}.$$

(These estimates are supposed to hold uniformly for r and R , but may depend on k .) Then we have for any $d \in \mathbb{N}$,

$$(2.8) \quad \sup_{x \in \mathbb{R}^d} \left| \int_{\mathbb{R}^d} \chi(|\xi|) e^{-i\varphi(|\xi|)t + i\xi x} d\xi \right| \lesssim t^{-d/2} (\varphi'(R)/R)^{-(d-1)/2} (\varphi''(R))^{-1/2},$$

where the constant depends only on φ, χ, d .

Proof. Since the case $d = 1$ is easy, we may assume $d \geq 2$. The phase can be stationary when its gradient

$$(2.9) \quad \Phi := -t\varphi'(r)\theta + x$$

vanishes, where we use the polar coordinates $\xi = r\theta$. Because all other parts are radially symmetric, we may assume that $x = |x|e_1$, where $e_1 = (1, 0, \dots, 0) \in \mathbb{R}^d$.

By the monotonicity of φ' , the stationary point for ξ is uniquely determined for any fixed (t, x) , by $\xi = ce_1$ where c is the unique solution (if it exists) for the equation

$$(2.10) \quad |x| = t\varphi'(c).$$

First we dispose of the region away from the stationary point. There exists a smooth partition of unity on S^{d-1} , denoted by $F_1(\theta), \dots, F_N(\theta) \in C^\infty(S^{d-1})$ where N depends only on d , satisfying the following.

1. $\sum_j F_j(\theta) = 1$ for any $\theta \in S^{d-1}$.
2. There exists $v^j \in S^{d-1}$ for each j such that $\theta \cdot v^j > 1/2$ for any $\theta \in \text{supp } F_j$.
3. For any $j > 1$, $v^j \cdot x < 0$.

Thus we have $\text{supp } F_j \subset \{\theta \cdot e_1 \leq \sqrt{3}/2\}$ for $j > 1$. Denote by I the integral to be estimated. We decompose the integral by using F_j and estimate those terms for $j > 1$ by partial integration in the v^j direction.

$$(2.11) \quad I = \sum_{j=1}^N I_j, \quad I_j = \int \chi(r) F_j(\theta) e^{-it\varphi(r) + i\xi \cdot x} d\xi$$

$$(2.12) \quad I_j = \int e^{ip} (\partial_j i p_j^{-1})^K (\chi(r) F_j(\theta)) d\xi,$$

for any $K \geq 0$, where we denote

$$(2.13) \quad p := -t\varphi(r) + \xi \cdot x, \quad \partial_j = v_j \cdot \nabla, \quad p_j = \partial_j p,$$

and $\partial_j i p_j^{-1}$ denotes the operator $f \mapsto \partial_j (i p_j^{-1} f)$, not the derivative of $i p_j^{-1}$. This convention applies to similar expressions of the form $\partial_* p_*^{-1}$ appearing later, but otherwise, we use ∂^α to denote the usual partial derivatives.

By the property of F_j and v^j , we have

$$(2.14) \quad p_j = v^j \cdot (-t\varphi'(r)\theta + x) < -t\varphi'(r)/2,$$

for any $\theta \in \text{supp } F_j$. Expanding the K -th power, we estimate

$$(2.15) \quad |I_j| \lesssim \sum_{|\alpha|=K} \int \left| \frac{\partial_j^{\alpha_1} p_j \cdots \partial_j^{\alpha_l} p_j}{p_j^{K+l}} \partial_j^{\alpha_0} (\chi(r) F_j(\theta)) \right| d\xi.$$

By using the estimates $|\varphi^{(k+1)}| \lesssim |\varphi'/r^k|$, $|\chi^{(k)}| \lesssim r^{-k}$ and $|\nabla^k F_j(\theta)| \lesssim r^{-k}$, we obtain

$$(2.16) \quad |I_j| \lesssim (t\varphi'(R)R)^{-K} R^d$$

for any $j > 1$ and $K \geq 0$. Choosing $K = d/2$, we obtain the desired estimate for those I_j , since $|\varphi''| \lesssim |\varphi'/r|$.

Now we are left with the stationary part I_1 . We further decompose this region into parallelepiped regions by a smooth partition of unity. Their side length should be $\delta := (t\varphi''(R))^{-1/2}$ in the e_1 direction and $\lambda := (t\varphi'(R)/R)^{-1/2}$ in the other directions. Denote the set of such parallelepipeds by \mathcal{P} . For each

$P \in \mathcal{P}$, we have a smooth cut-off function $G_P(\xi)$. We have $\sum_{P \in \mathcal{P}} G_P(\xi) = 1$ on $\text{supp } F_1(\theta)\chi(r)$, $\xi_1 \sim |\xi| \sim R$ on any $P \in \mathcal{P}$, and

$$(2.17) \quad |D^\alpha G_P| \lesssim \delta^{-\alpha_1} \lambda^{-\alpha_2 - \dots - \alpha_n}.$$

Denote the integral corresponding to G_P by I_P .

We put the stationary point at the center of one parallelepiped P_0 . In that region, the integral is not very oscillatory, so we just bound it by the volume:

$$(2.18) \quad |I_{P_0}| \lesssim \delta \lambda^{d-1},$$

which yields the desired estimate for this part.

We have the same volume bound on the other parts, but we need additional decay factors to sum those numerous pieces up. For that purpose we may integrate by parts either in the e_1 direction or in the other $e^1 \perp e_1$ directions.

First we consider integration in x^1 direction. Take any $P \in \mathcal{P}$ on which we have $|\xi^1| \sim \lambda k$, $k \in \mathbb{N}$. Notice that such a P exists only if $\lambda \lesssim R$. Then we have a direction $v \perp e_1$ such that $v \cdot \xi \sim \lambda k$ for any $\xi \in P$. We integrate by parts in the v direction. Then we obtain

$$(2.19) \quad I_P = \int e^{-it\varphi(r) + i\xi \cdot x} (\partial_v i p_v^{-1})^K (\chi(r) F_1(\theta) G_P(\xi)) d\xi,$$

for any $K \geq 0$, where $\partial_v = v \cdot \nabla$ and

$$(2.20) \quad p_v = \partial_v(-t\varphi(r) + \xi \cdot x) = -tv \cdot \xi \varphi'(r)/r \sim -t\lambda k \varphi'(R)/R.$$

By expanding the K -th power, we obtain

$$(2.21) \quad |I_P| \lesssim (t\lambda k \varphi'(R) R \min(\lambda, R))^{-K} \lambda^{d-1} \delta,$$

where λ in $\min(\lambda, R)$ is coming from derivatives hitting $G_P(\xi)$. Thus we get

$$(2.22) \quad |I_P| \lesssim k^{-K} \delta \lambda^{d-1}.$$

Next we consider integration in the e_1 direction. Here we should be a bit more careful, since the phase can be very small depending on φ'' even if we are not in P_0 , and then the estimate on the higher derivatives of the phase becomes less trivial. Take any $P \in \mathcal{P}$ on which we have $\xi_1 - c \sim \delta j$ with $j \neq 0$, where c denotes the ξ_1 value of the stationary point. Such a P exists only if $\delta \lesssim R$. Let $p_1 = \partial_1 p = -t\varphi'(r)\theta_1 + |x|$. We compute

$$(2.23) \quad \begin{aligned} \partial_1 r &= \theta_1, & \partial_1 \theta_1 &= (1 - \theta_1^2)/r, \\ \partial_1 p_1 &= -t(\varphi''\theta_1 + \varphi'(1 - \theta_1^2)/r) \end{aligned}$$

For higher derivatives we have

$$(2.24) \quad |\partial_1^{k+1} p_1| \lesssim \sum_{j=0}^k t |\varphi^{(2+j)}/r^j| + t |\varphi'(1 - \theta_1^2)/r^{k+1}|.$$

Since $\varphi', \varphi'' > 0$, we have $|\partial_1^{k+1} p_1| \lesssim |\partial_1 p_1|/r^k$. This in particular implies that $\partial_1 p_1$ is of the same size on $\text{supp } \chi$. Then it implies that $|p_1| \sim |\partial_1 p_1 \delta j|$. We have also $|\partial_1 p_1| \gtrsim t|\varphi''(R)|$.

Integrating by parts in the e_1 direction we have

$$(2.25) \quad I_P = \int e^{-it\varphi(r)+i\xi \cdot x} (\partial_1 i p_1^{-1})^K (\chi(r) F_1(\theta) G_P(\xi)) d\xi ,$$

Expanding the K -th power, we estimate

$$(2.26) \quad |I_P| \lesssim (t\varphi''(R)\delta j\delta)^{-K} \lambda^{d-1} \delta,$$

Thus we obtain

$$(2.27) \quad |I_P| \lesssim j^{-K} \lambda^{d-1} \delta.$$

In conclusion, we have for any $P \in \mathcal{P}$,

$$(2.28) \quad |I_P| \lesssim (1+k+j)^{-K} \lambda^{d-1} \delta,$$

for any $K \in \mathbb{N}$, so we can sum them up for all k, j , deriving the desired estimate. \square

Proof of Theorem 2.1. $\varphi(r) = r\sqrt{r^2 + 2}$ satisfies the assumption of the above theorem and

$$(2.29) \quad \varphi'(r) \sim r + 1, \quad \varphi''(r) \sim \frac{r}{r + 1}.$$

Thus (2.8) gives the desired L^∞ decay estimate, i.e., (2.5) for $q = \infty$. It is extended to $q < \infty$ by interpolation with the trivial L^2 bound $q = 2$. Then we derive the Strichartz estimate (2.6) from (2.5), see [19]. \square

3. Low frequency normal form

Let u be a solution of (1.5) and $v = V^{-1}u$. By using (2.4), we can derive the equation for v :

$$(3.1) \quad v(t) = e^{-itH} v_0 - \int_0^t e^{-i(t-s)H} V^{-1} iF(u(s)) ds.$$

The presence of U^{-1} in the nonlinearity is our only “enemy” for $d \geq 4$; without it, the scattering would follow in the same way as for NLS with quadratic and cubic power nonlinearities. Quadratic nonlinearities have the critical decay rate for scattering by the Strichartz estimate when $d = 4$, while for $d > 4$ their decay is more than sufficient. Indeed we can derive the scattering result for $d > 4$ and certain restricted values of σ by using only the Strichartz estimate with the U gain and the Sobolev embedding. That argument, however, cannot work in the critical case $d = 4$.

We introduce the following nonlinear change of variable to kill the singular nonlinearity. It gives a better range of σ even for $d > 4$. Let

$$(3.2) \quad w := N(u) := u + P|u|^2/2,$$

where P is the Fourier multiplier $\mathcal{F}^{-1}\chi\mathcal{F}$, $\chi \in C_0^\infty(\mathbb{R}^d)$ satisfying $\chi(x) = 1$ for $|x| \leq 1$ and $\chi(x) = 0$ for $|x| \geq 2$. Then w satisfies the following equation:

$$(3.3) \quad \begin{aligned} i\dot{w} &= -\Delta w + 2\operatorname{Re} w + G_1 + iG_2, \\ G_1 &= (3 - P)u_1^2 + Qu_2^2 + P\Delta|u|^2/2 + |u|^2u_1, \\ G_2 &= 2Q(u_1u_2) + \nabla P \cdot (u_2\nabla u_1 - u_1\nabla u_2) + Q(|u|^2u_2), \end{aligned}$$

where $u = u_1 + iu_2$ and $Q = Id - P$. Thus

$z := V^{-1}w = V^{-1}N(u) = U^{-1}(u_1 + P|u|^2/2) + iu_2 = v + U^{-1}P|Vv|^2/2 =: Mv$ satisfies

$$(3.4) \quad z = e^{-iHt}z(0) - \int_0^t e^{-iH(t-s)}(iG_1 - U^{-1}G_2)ds.$$

Notice that the singularity of U^{-1} in front of G_2 is now canceled by either Q or ∇ . There is no loss of regularity thanks to P . Moreover, G does not contain bilinear interaction of very low frequency of u_2 , whereas u_1 behaves better at the Fourier 0 due to the relation with z .

4. Proof of the main theorem

We introduce Besov spaces with different regularity for low and high frequency:

$$(4.1) \quad \|\varphi\|_{B_q^{a,b}} := \|P\varphi\|_{\dot{B}_q^a} + \|Q\varphi\|_{\dot{B}_q^b}, \quad H^{a,b} := B_2^{a,b},$$

where $\dot{B}_q^a := \dot{B}_{q,2}^a$ denotes the homogeneous Besov space (cf. [1]). We will use these spaces under the restriction $a < d/q$, embedding them into \mathcal{S}' . Moreover, we will mainly use them in the case $a > d/q - d$, where C_0^∞ is a dense subspace. We have $B_q^{a,b} = U^{b-a}\dot{B}_{q,2}^b$, $H^{0,b} = H^b$ and the embedding $B_q^{a,b} \subset B_q^{a',b'}$ for $a \leq a'$ and $b \geq b'$. These spaces have the following multiplicative property, which is enough for us to estimate the nonlinearity.

Lemma 4.1. *Let $p_1, p_2, p_3 \in [2, \infty]$, $s_1, s_2, s_3 \in \mathbb{R}$, $t_1, t_2, t_3 \in \mathbb{R}$ and $b_a := 1/p_a$ for $a = 1, 2, 3$. Assume that for $a = 1, 2, 3$,*

$$\max(0, s_a, s_1 + s_2 + s_3) \leq d(b_1 + b_2 + b_3 - 1) \leq t_1 + t_2 + t_3 \geq t_a, \quad s_a < db_a.$$

Then we have

$$(4.2) \quad \left| \int_{\mathbb{R}^d} fgh dx \right| \lesssim \|f\|_{B_{p_1}^{s_1, t_1}} \|g\|_{B_{p_2}^{s_2, t_2}} \|h\|_{B_{p_3}^{s_3, t_3}}.$$

Proof. We apply the Littlewood-Paley decomposition to each of the functions.

$$(4.3) \quad \int fgh = \sum_{l=-\infty}^{\infty} \sum_{k=l-2}^l \sum_{j=-\infty}^k c_{j,k,l} \int f_j g_k h_l + f_k g_l h_j + f_l g_j h_k,$$

where $0 \leq c_{j,k,l} \leq 1$ and each f_j is supported in Fourier space in $\{2^{j-1} \leq |\xi| \leq 2^{j+1}\}$. By symmetry, it suffices to estimate the first summand. Let

$\beta = 1 - b_2 - b_3$, then we have $0 \leq \beta \leq b_1$. By the Sobolev and the Hölder inequalities and the embedding $\dot{B}_p^0 \subset L^p$, we have

$$(4.4) \quad \left| \sum_{j \leq k} \int f_j g_k h_l \right| \leq \left\| \sum_{j \leq k} f_j \right\|_{L^{1/\beta}} \|g_k\|_{L^{p_2}} \|h_l\|_{L^{p_3}} \\ \lesssim \left\| \sum_{j \leq k} 2^{j(-s+d[b_1-\beta])} f_j \right\|_{\dot{B}_{p_1}^s} \|g_k\|_{L^{p_2}} \|h_l\|_{L^{p_3}}$$

for any s . Taking $s = s_1$ for $j \leq 0$ and $s = t_1$ for $j > 0$ (when $k > 0$), the norm on f is bounded by

$$(4.5) \quad \begin{cases} 2^{k(-s_1+d(b_1-\beta))} & (k \leq 0), \\ 1 + 2^{k(-t_1+d(b_1-\beta))} & (k > 0), \end{cases}$$

since $-s_1 + d(b_1 - \beta) \geq 0$. The summation over $k \sim l \in \mathbb{Z}$ converges, because for $k \leq 0$,

$$(4.6) \quad -s_1 + d(b_1 - \beta) - s_2 - s_3 = -(s_1 + s_2 + s_3) + d(b_1 + b_2 + b_3 - 1) \geq 0,$$

and for $k \geq 0$,

$$(4.7) \quad -t_2 - t_3 \leq 0, \quad -t_1 - t_2 - t_3 + d(b_1 + b_2 + b_3 - 1) \leq 0,$$

and the summability is provided by the ℓ^2 norm in the Besov norms. \square

For the proof of our main theorem, we will use the following special cases of the above lemma. Let $d \geq 4$, $s = d/2 - 1$, $b = 1/p = 1 - 1/p' = 1/2 - 1/d$ and $1/q = 1/2 - 1/(2d)$. Then we have

$$(4.8) \quad \begin{aligned} (1) \quad & H^{\sigma-b, s-j} \times B_p^{\sigma-b, s-k} \subset B_{p'}^{\sigma+b, s-j-k}, \\ (2) \quad & H^{\sigma, s} \times B_p^{\sigma-b, s-1} \subset H^{\sigma-b, s-1}, \\ (3) \quad & B_q^{\sigma-b/2, s} \times B_q^{\sigma-b/2, s} \subset H^{\sigma, s}, \\ (4) \quad & H^{\sigma, s} \times H^{\sigma, s} \subset B_{p'}^{\sigma, s}, \\ (5) \quad & H^{\sigma, s} \times B_q^{\sigma-b/2, s} \subset H^{\sigma-1, s-1/2}, \end{aligned}$$

for $0 \leq j, k, j+k \leq 1$, provided that $|\sigma| - b \leq d/2 - 2$. All conditions in the lemma are satisfied because

$$(4.9) \quad \begin{aligned} (1) \quad & \max(\sigma - 3b, |\sigma| - b) \leq d/2 - 2 \leq s \geq j + k - s, \quad |\sigma| - b < d/2 - 1, \\ (2) \quad & \max(\sigma, b - \sigma) \leq d/2 - 1 \leq s \geq 1 - s, \quad \sigma < d/2, \\ (3) \quad & \max(\sigma - b/2, -\sigma) \leq d/2 - 1 \leq s, \quad \sigma - b/2 < d/2 - 1/2, \\ (4) \quad & |\sigma| \leq d/2 - 1 \leq s, \\ (5) \quad & \max(\sigma - b/2 + 1, \sigma, 1 - \sigma) \leq d/2 - 1/2 \leq s + 1/2 \geq 1/2 - s. \end{aligned}$$

As the function space for $v = V^{-1}u$ and $z = V^{-1}N(u)$, we define X_0 and X_1 by

$$(4.10) \quad X_j = L^\infty H^{\sigma-jb, s} \cap L^2 B_p^{\sigma-b, s}.$$

We have $U : X_0 \rightarrow X_1 \subset X_0$ bounded. First we consider the norm relation between v and z . By (4.8)-(2)(4) and the embeddings $B_{p'}^{\sigma,s} \subset H^{\sigma-1,s-1}$ and $H^{\sigma-b,s-1} \subset B_p^{\sigma-b-1,s-2}$, we have

$$(4.11) \quad \|U^{-1}P(fg)\|_B \lesssim \|f\|_{H^{\sigma,s}} \|g\|_B$$

for $B = H^{\sigma,s}$ and $B = B_p^{\sigma-b,s}$. Hence for any small $z \in H^{\sigma,s}$, the unique inverse $v = M^{-1}z$ is obtained in the same space by applying the Banach fixed point theorem to the equation $v = z - U^{-1}P|Vv|^2/2$. Thus M is bi-Lipschitzian for small data in $H^{\sigma,s}$. In particular, we have the following estimate for the free part $z_0 := e^{-iHt}z(0)$

$$(4.12) \quad \|z_0\|_{X_0} \lesssim \|z(0)\|_{H^{\sigma,s}} \lesssim \|v(0)\|_{H^{\sigma,s}} < \delta.$$

Moreover, (4.11) implies that $\|z\|_{X_0} \sim \|v\|_{X_0}$ as long as they remain small. In addition, we have

$$(4.13) \quad \|M(e^{-iHt}\varphi) - e^{-iHt}\varphi\|_{H^{\sigma,s}} \rightarrow 0 \quad (|t| \rightarrow \infty)$$

for any $\varphi \in H^{\sigma,s}$. This is proved as follows. By our assumption $|\sigma| \geq -d/2+2-b$, C_0^∞ is a dense subspace of $H^{\sigma,s}$ and $B_{q'}^{\sigma+b/2,s}$. The decay estimate (2.5) and (4.8)-(5) imply that for any $\varphi \in C_0^\infty$,

$$(4.14) \quad \begin{aligned} \|U^{-1}P|Ve^{-iHt}\varphi|^2\|_{H^{\sigma,s}} &\lesssim \|e^{-iHt}\varphi\|_{H^{\sigma,s}} \|e^{-iHt}\varphi\|_{B_q^{\sigma-b/2,s}} \\ &\lesssim |t|^{-1/2} \|\varphi\|_{H^{\sigma,s}} \|\varphi\|_{B_{q'}^{\sigma+b/2,s}}. \end{aligned}$$

Hence we obtain (4.13) first for $\varphi \in C_0^\infty$, and then extend it to general $\varphi \in H^{\sigma,s}$ by the denseness and the local Lipschitz continuity of M . The bi-Lipschitz property implies also that

$$(4.15) \quad \|e^{-iHt}\varphi - M^{-1}(e^{-iHt}\varphi)\|_{H^{\sigma,s}} \rightarrow 0 \quad (|t| \rightarrow \infty)$$

for any small $\varphi \in H^{\sigma,s}$.

Next we estimate the quadratic terms in $G(u)$. By (4.8)-(1), we have

$$(4.16) \quad \begin{aligned} \|u_1u\|_{L^2B_{p'}^{\sigma+b,s}} &\lesssim \|u_1\|_{L^\infty H^{\sigma-b,s}} \|u\|_{L^2B_p^{\sigma-b,s}} \lesssim \|u_1\|_{X_1} \|u\|_{X_0}, \\ \|P(u\nabla u)\|_{L^2B_{p'}^{\sigma+b,s}} &\lesssim \|u\|_{L^2B_p^{\sigma-b,s}} \|\nabla u\|_{L^\infty H^{\sigma-b,s-1}} \lesssim \|u\|_{X_0}^2, \end{aligned}$$

These are used to estimate the terms u_1^2 , u_1u_2 , $P\Delta|u|^2$ and $u_2\nabla u_1 - u_1\nabla u_2$, together with the boundedness of $U^{-1}Q$, $U^{-1}\nabla P$ and $P\nabla$ on every Besov space. Qu_2^2 can be treated in the same way as u_1u since

$$(4.17) \quad Qu_2^2 = Q(u_2 \cdot Q_{-2}u_2) + Q(Q_{-2}u_2 \cdot P_{-2}u_2)$$

and $Q_{-2} : X_0 \rightarrow X_1$ bounded, where $P_{-2} := \mathcal{F}^{-1}\chi(4\xi)\mathcal{F}$ and $Q_{-2} := Id - P_{-2}$, defined in the same way as P and Q (see below (3.2)).

For the cubic terms, we use (4.8)-(3)(4) together with the interpolation $X_0 \subset L^4B_q^{\sigma-b/2,s}$, yielding

$$(4.18) \quad \|u^3\|_{L^2B_{p'}^{\sigma,s}} \lesssim \|u\|_{L^\infty H^{\sigma,s}} \|u^2\|_{L^2H^{\sigma,s}} \lesssim \|u\|_{X_0} \|u\|_{L^4B_q^{\sigma-b/2,s}}^2 \leq \|u\|_{X_0}^3.$$

Gathering these estimates, we obtain

$$(4.19) \quad \begin{aligned} \|z - z_0\|_{X_0} &\lesssim \|V^{-1}iG\|_{L^2 B_{p'}^{\sigma+b,s}} \\ &\lesssim \|u\|_{X_0} (\|u\|_{X_0} + \|u_1\|_{X_1} + \|u\|_{X_0}^2) \lesssim \|z\|_{X_0}^2, \end{aligned}$$

if the right hand side is small. By continuity in time and (4.12), we conclude that

$$(4.20) \quad \|z\|_{X_0} + \|z_0\|_{X_0} + \|V^{-1}iG\|_{L^2 B_{p'}^{\sigma+b,s}} \lesssim \|v(0)\|_{H^{\sigma,s}}.$$

In particular, the norm of G on the interval (T, ∞) vanishes as $T \rightarrow \infty$, which implies via the Strichartz estimate that $e^{iHt}z(t) \rightarrow \exists v_{\pm}$ strongly in $H^{\sigma,s}$. Then (4.15) implies that $e^{iHt}v(t) \rightarrow v_{\pm}$ in $H^{\sigma,s}$.

We can construct the wave operators applying the same estimates to the Cauchy problem from $t = \pm\infty$:

$$(4.21) \quad z = e^{-iHt}v_{\pm} - \int_{\pm\infty}^t e^{-iH(t-s)}V^{-1}iG(VM^{-1}z(s))ds.$$

The bi-Lipschitz property of the maps $v(0) \mapsto v_{\pm}$ is similarly proved by estimating the difference of two solutions in the same space. These arguments, essentially by the Banach fixed point theorem, are quite standard in nonlinear scattering theory, and so we omit the details.

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