LECTURE 2

Mathematical tools from potential theory: capacities, harmonic functions, variational principles.

\S SETTING

We place ourselves in the setting of a discrete-time Markov process $X = (X_n)_{n \in \mathbb{N}_0}$ on a countable state space S with a transition kernel

$$P = (p(x, y)_{x, y \in S})$$

and a generator $L = P - \mathbb{I}$. We assume that X is irreducible.

1. The first fundamental object in potential theory is the following. Let $A, B \subset S$ be two non-empty disjoint subsets. Consider the Dirichlet problem

$$(-Lh)(x) = 0, \quad \forall x \in S \setminus (A \cup B),$$

 $h(x) = 1, \quad \forall x \in A,$
 $h(x) = 0, \quad \forall x \in B.$



Dirichlet problem for $h: S \to [0, 1]$ with boundary conditions $h \equiv 1$ on A ('target set') and $h \equiv 0$ on B ('killing set').

If S is finite, then the Dirichlet problem always has a unique solution. This solution, which is harmonic on $S \setminus (A \cup B)$, is denoted by $h_{A,B}(x)$ and is called the equilibrium potential.

2. The following representation holds:

$$h_{A,B}(x) = \mathbb{P}_x (\tau_A < \tau_B), \quad x \in S \setminus (A \cup B).$$

Here, $\tau_C = \inf\{n \in \mathbb{N} : X_n \in C\}$, where X_0 is ignored.

REMARK: For $x \in A \cup B$, we can write

$$\mathbb{P}_x \left(\tau_A < \tau_B \right) = \sum_{y \in S \setminus (A \cup B)} p(x, y) \mathbb{P}_y \left(\tau_A < \tau_B \right) + \sum_{y \in A} p(x, y)$$
$$= \sum_{y \in S} p(x, y) h_{A,B}(y) = (Ph_{A,B})(x)$$
$$= (Lh_{A,B})(x) + h_{A,B}(x).$$

Hence

$$x \in A: \quad (-Lh_{A,B})(x) = \mathbb{P}_x (\tau_B < \tau_A),$$

$$x \in B: \quad (Lh_{A,B})(x) = \mathbb{P}_x (\tau_A < \tau_B).$$

The quantity

$$e_{A,B}(x) \equiv (-Lh_{A,B})(x), \qquad x \in A,$$

is called the equilibrium measure on A, and is the second fundamental object. We know that $e_{A,B}(x) = 0$ for all $x \in S \setminus (A \cup B)$.

3. The Green function of X killed at B is defined as

$$G_B(x,y) = \mathbb{E}_x \left[\sum_{n=0}^{\tau_B - 1} \mathbf{1}_{\{X_n = y\}} \right], \quad x, y \notin B,$$

and is the third fundamental object.

4. The equilibrium potential $h_{A,B}$ satisfies the following inhomogeneous Dirichlet problem

$$(-Lh)(x) = e_{A,B}(x), \quad \forall x \in S \setminus B,$$

 $h(x) = 0, \quad \forall x \in B.$

THEOREM 2.1

$$h_{A,B}(x) = \sum_{y \in A} G_B(x, y) e_{A,B}(y), \qquad x \in S.$$

THEOREM 2.1, whose proof is straightforward, allows us to express the Green function in terms of the equilibrium potential and the equilibrium measure: simply choose $A = \{a\}$, to get

$$G_B(x,a) = \frac{h_{a,B}(x)}{e_{a,B}(a)}, \qquad x \in S.$$

Note that

$$e_{a,B}(a) = \mathbb{P}_a(\tau_B < \tau_a)$$

has the meaning of an escape probability from $a \in A$ to B.

§ DIRICHLET FORM

Henceforth we restrict to reversible Markov processes.

DEFINITION 2.2

A Markov process with countable state space S and with transition kernel $P = (p(x, y))_{x,y \in S}$ is called reversible if there exists $\mu: S \to [0, \infty)$ such that

$$\mu(x)p(x,y) = \mu(y)p(y,x) \qquad \forall x, y \in S.$$

The function μ is called the reversible measure of the Markov process, and is easy to compute except for the normalisation.

The function space $L^2(S,\mu)$ is a natural space to work on when the Markov process is reversible with respect to μ .

LEMMA 2.3

(a) If μ is a reversible measure for P, then μ is an invariant measure for P.

(b) If μ is an invariant measure for P, then $Pf \in L^2(S,\mu)$ for all $f \in L^2(S,\mu)$.

DEFINITION 2.4

Let L be the generator of a Markov process with reversible measure μ . Then L defines a quadratic form

$$\mathcal{E}(f,g) \equiv \sum_{x \in S} \mu(x) f(x) (-Lg)(x), \qquad f,g \in L^2(S,\mu),$$

called the Dirichlet form, which is non-negative-definite.

In the discrete case it is easy to write out $\mathcal{E}(f,g)$ explicitly. Namely, by reversibility,

$$\mathcal{E}(f,g) = \sum_{x,y\in S} \mu(x)p(x,y)f(x)[g(x) - g(y)]$$
$$= \sum_{x,y\in S} \mu(x)p(x,y)f(y)[g(y) - g(x)].$$

Symmetrising between the first and the second expression, we get

$$\mathcal{E}(f,g) = \frac{1}{2} \sum_{x,y \in S} \mu(x)p(x,y) \\ \times \left\{ f(x)[g(x) - g(y)] + f(y)[g(y) - g(x)] \right\} \\ = \frac{1}{2} \sum_{x,y \in S} \mu(x)p(x,y) \\ \times [f(x) - f(y)][g(x) - g(y)].$$

THEOREM 2.5

If P is reversible with respect to μ , then for all non-empty disjoint sets $A, B \subset S$,

$$h_{A,B}(x) = \sum_{y \in A} \frac{\mu(y)}{\mu(x)} G_B(y, x) e_{A,B}(y), \qquad x \in S.$$

In particular, if f given g is a solution of the inhomogeneous Dirichlet problem

$$\begin{array}{rcl} (-Lf)(x) &=& g(x), & \forall x \in S \setminus B, \\ f(x) &=& 0, & \forall x \in B, \end{array}$$

then

$$\sum_{y \in A} \nu_{A,B}(y) f(y) = \frac{1}{\operatorname{cap}(A,B)} \sum_{x \in S} \mu(x) h_{A,B}(x) g(x),$$

where $\nu_{A,B}$ is the probability measure on A given by

$$\nu_{A,B}(y) = \frac{\mu(y)e_{A,B}(y)}{\operatorname{cap}(A,B)}, \qquad y \in A,$$
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with normalisation factor the capacity

$$\operatorname{cap}(A,B) \equiv \sum_{x \in A} \mu(x) e_{A,B}(x).$$

PROOF:

By reversibility,

$$\mu(x)G_B(x,y) = \mu(y)G_B(y,x), \qquad x, y \in S,$$

which yields the formula for $h_{A,B}$ via THEOREM 2.1.

Multiplying this formula by $\mu(x)g(x)$, summing over $x \in S$, and noting that $\sum_{x \in S} G_B(y, x)g(x) = f(y)$, we get the formula for $\sum_{y \in A} \nu_{A,B}(y)f(y)$, apart from the normalisation factor cap(A, B).

Here, $\nu_{A,B}$ is called the last-exit biased distribution on A for the transition from A to B, while cap(A,B) is called the capacity of the pair (A,B).

The following corollary of THEOREM 2.5, obtained by picking g = 1 on $S \setminus B$, offers a formula for mean hitting times that plays a crucial role in our study of metastability.

COROLLARY 2.6

Let $A, B \subset S$ be non-empty and disjoint. For reversible Markov processes,

$$\sum_{x \in A} \nu_{A,B}(x) \mathbb{E}_x[\tau_B] = \frac{1}{\operatorname{cap}(A,B)} \sum_{y \in S} \mu(y) h_{A,B}(y).$$

In particular, for $A = \{x\}$,

$$\mathbb{E}_x[\tau_B] = \frac{1}{\operatorname{cap}(x,B)} \sum_{y \in S} \mu(y) h_{x,B}(y).$$

LEMMA 2.7

Let $A, B \subset S$ be non-empty and disjoint. Then cap(A, B) can be expressed as

$$\operatorname{cap}(A,B) = \mathcal{E}(h_{A,B},h_{A,B}).$$

PROOF:

This is obvious from the definition of the Dirichlet form, the equilibrium potential, the equilibrium measure, and the capacity.

Note that LEMMA 2.7 becomes particularly useful through the alternative representation of the Dirichlet form as a quadratic form.

\S DIRICHLET PRINCIPLE



The capacity of A, B is given by the Dirichlet Principle

$$\operatorname{cap}(A,B) = \inf_{\phi \in \Phi_{A,B}} \mathcal{E}(\phi,\phi)$$

where

$$\Phi_{A,B} = \{\phi \colon S \to [0,1] \colon \phi(A) = 1, \, \phi(B) = 0\}$$

and

$$\mathcal{E}(\phi,\phi) = \sum_{x,y\in S} \mu(x)p(x,y)[\phi(y) - \phi(x)]^2$$

is the Dirichlet form associated with the Markov dynamics.

\S THOMSON PRINCIPLE



A unit flow from B to A is a map $u: S \times S \to \mathbb{R}$ such that the flows into and out of nodes in $S \setminus \{A, B\}$ equal 0, while the flow out of B and into A equal 1.

The Thomson Principle reads

$$\mathsf{cap}(A,B) = \sup_{u \in \mathcal{U}_{A,B}} \frac{1}{\mathcal{D}(u,u)}$$

where $\mathcal{U}_{A,B}$ is the set of unit flows from B to A, and

$$\mathcal{D}(u,u) = \sum_{x,y \in S} \frac{1}{\mu(x)p(x,y)} u(x,y)^2$$

is a dual of the Dirichlet form.

The infimum in the Dirichlet Principle is uniquely taken at the equilibrium potential

$$\phi(x) = h_{A,B}(x),$$

while the supremum in the Thomson Principle is uniquely taken at the equilibrium flow u given by

$$u(x,y) = \frac{\mu(x)p(x,y)\left[h_{A,B}(y) - h_{A,B}(x)\right]_{+}}{\operatorname{cap}(A,B)}.$$

\S CAPACITY ESTIMATES

The estimation of capacity proceeds via

- Dirichlet principle $\operatorname{cap}(A,B) \leq \mathcal{E}(\phi,\phi),$
- Thomson principle $\operatorname{cap}(A,B) \geq 1/\mathcal{D}(u,u)$,

where ϕ, u are properly chosen test functions and test flows that live in the vicinity of the critical droplet.

The choice of ϕ , u requires physical insight into what drives the metastable crossover. 17





The formula relating metastable crossover time to capacity effectively links non-equilibrium to equilibrium. The inverse of the capacity plays the role of effective resistance.

ASYMPTOTICS:

The Dirichlet Principle and Thomson Principle allow for the derivation of upper and lower bounds on capacity. With care, these can be made to match asymptotically.

REDUCTION:

In metastable regimes the high-dimensional Dirichlet form and dual Dirichlet form are largely controlled by the low-dimensional set of critical droplets.



§ RANDOM WALKS

For later use, in Lecture 4, we compute relevant quantities for nearest-neighbour random walk on an interval $[a, b] \subset \mathbb{Z}$.

• The Dirichlet problem reads

$$p(x, x + 1)[h(x + 1) - h(x)] + p(x, x - 1)[h(x - 1) - h(x)] = 0, \quad a < x < b,$$

$$h(a) = 0,$$

$$h(b) = 1.$$

This is a recursion relation for the differences h(x) - h(x - 1), and a straightforward computation leads to expression (swap the roles of a and b)

$$h_{b,a}(x) = \mathbb{P}_x(\tau_b < \tau_a) = \frac{R(a, x)}{R(a, b)}, \qquad a < x < b,$$

with (use reversibility to bring in μ)

$$R(u,v) = \sum_{y=u+1}^{v} \frac{1}{\mu(y)} \frac{1}{p(y,y-1)}, \qquad u < v.$$

• The equilibrium measure is given by the formula

$$e_{a,b}(a) = p(a, a + 1)h_{b,a}(a + 1) + p(a, a - 1)h_{b,a}(a - 1) = p(a, a + 1)h_{b,a}(a + 1),$$

where we use that $h_{b,a}(a-1) = 0$. Inserting the formula for $h_{b,a}$ derived above, we get

$$e_{a,b}(a) = p(a, a + 1) \frac{R(a, a + 1)}{R(a, b)}$$
$$= \frac{\frac{p(a, a + 1)}{\mu(a + 1)p(a + 1, a)}}{R(a, b)} = \frac{1}{\mu(a)R(a, b)}.$$

Consequently, for the capacity we get

$$\operatorname{cap}(a,b) = rac{1}{R(a,b)}.$$

• Inserting the expressions for $h_{a,b}$ and cap(a,b) into the formula for $\mathbb{E}_x[\tau_a]$ in COROLLARY 2.6 (with $A = \{x\}$ and $B = \{a\}$), we get

$$\mathbb{E}_{x}[\tau_{a}] = R(a, x) \left(\sum_{y=a+1}^{x-1} \mu(y) \frac{R(a, y)}{R(a, x)} + \sum_{y=x}^{\infty} \mu(y) \right), \quad a < x.$$

This formula will be needed in Lecture 4 to compute the average metastable crossover time for a mean-field model called the Curie-Weiss model. The latter will be shown to link up with the Kramers formula for Brownian motion in a double-well potential, hinted at in Lecture 1.

\S dicrete versus continuous

Definitions and computations become more involved when the state space is infinite discrete or continuous.

Often A and B are not single configurations but are sets of configurations with an interesting geometric structure.



We will see examples in Lectures 9-12.

LITERATURE:

Chapter 7 of Bovier and den Hollander 2015, and references therein.